

Spin Transport in Kitaev Quantum Spin Liquid

Shi Feng and Nandini Trivedi

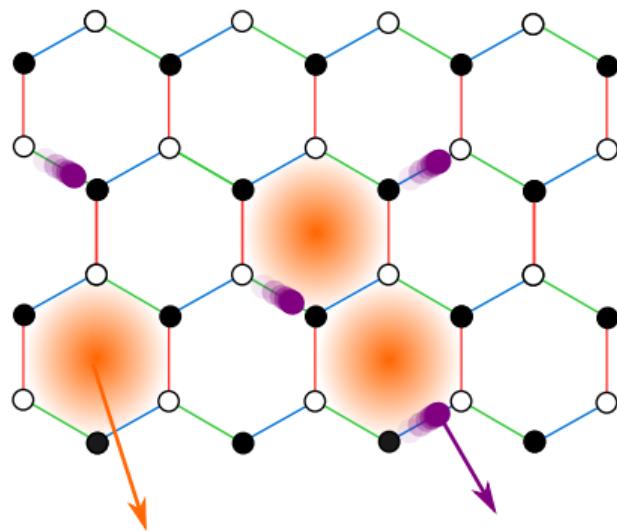
Department of Physics, The Ohio State University, Columbus, OH



Introduction

- A. Kitaev. Ann. Phys. 321, 2–111 (2006)

$$\mathcal{H} = \sum_{\alpha} \sum_{\langle ij \rangle} \textcolor{blue}{K_x} S_i^x S_j^x + \textcolor{green}{K_y} S_i^y S_j^y + \textcolor{red}{K_z} S_i^z S_j^z$$



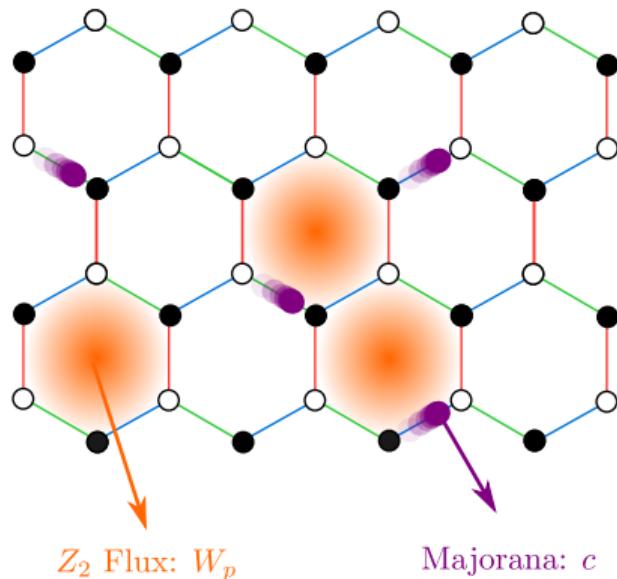
Z_2 Flux: W_p

Majorana: c

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Relevant QSL materials:

- α -RuCl₃

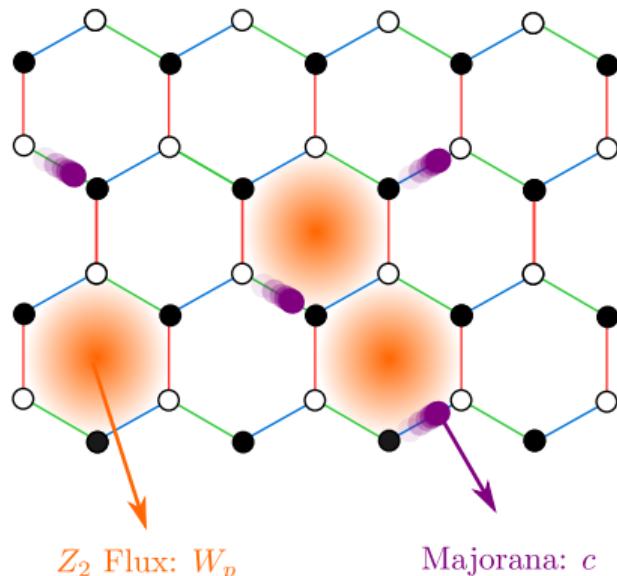
- ① Y. Kasahara, T. Ohnishi, Y. Mizukami, O. Tanaka, Sixiao Ma, K. Sugii, N. Kurita, H. Tanaka, J. Nasu, Y. Motome, T. Shibauchi and Y. Matsuda. Nature 559, 227–231 (2018)

Large in-plane field to destabilize order

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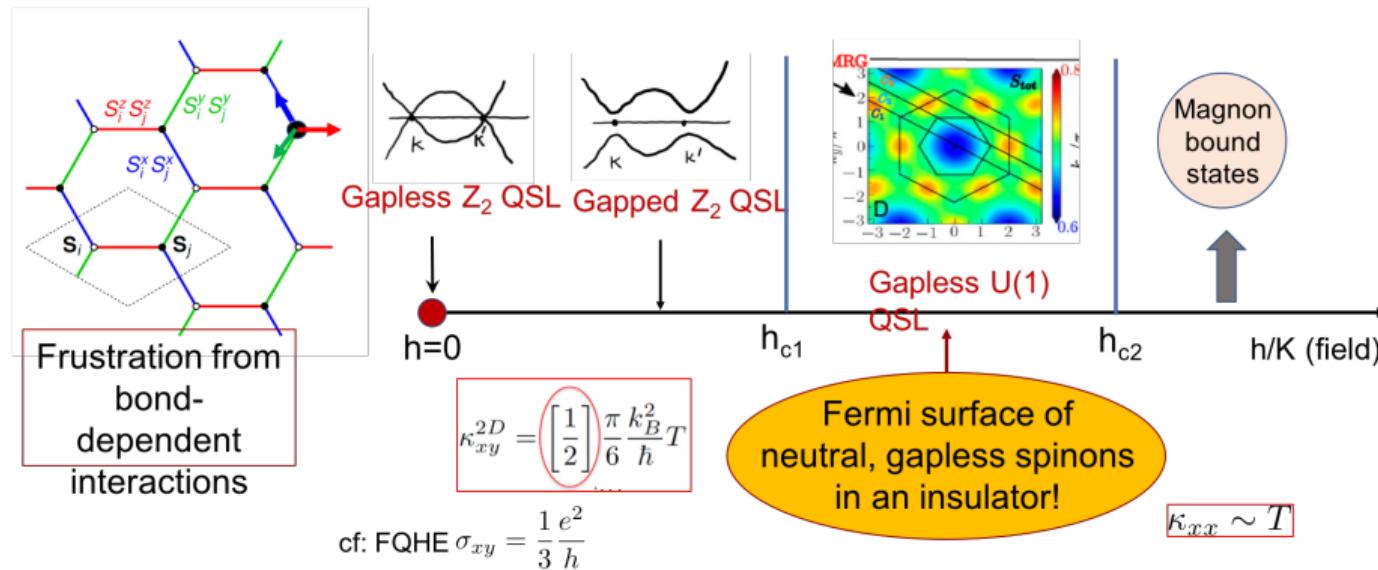
Large in-plane field to destabilize order

- BaCo₂(AsO₄)₂

① Xinshu Zhang, Yuanyuan Xu, Ruidan Zhong, R. J. Cava, N. Drichko, N. P. Armitage. In- and out-of-plane field induced quantum spin-liquid states in a more ideal Kitaev material: BaCo₂(AsO₄)₂. arXiv:2106.13418 (2021)

Small [111] field to destabilize order

Phase diagram of Kitaev model under [111] magnetic field

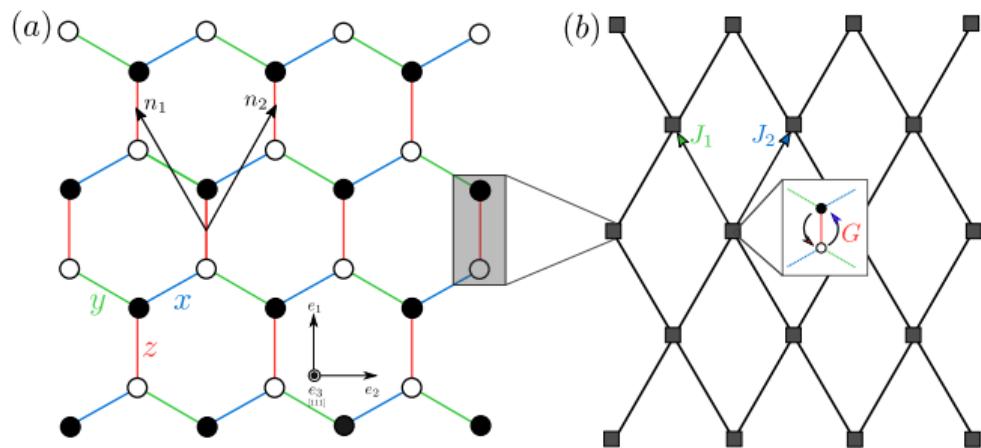


Niravkumar D. Patel and Nandini Trivedi. Magnetic field-induced intermediate quantum spin liquid with a spinon Fermi surface. PNAS, 116 (25) 12199-12203

Spin Transport in response to [111] field?

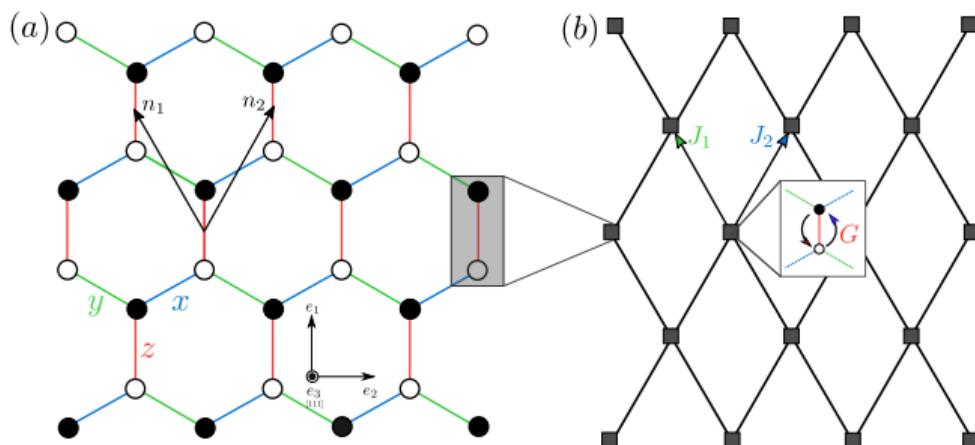
Spin transport in Kitaev model

$$\mathcal{H} = \sum_{\alpha} \sum_{\langle ij \rangle} K_{ij}^{\alpha} S_i^{\alpha} S_j^{\alpha} + \sum_i \frac{h^{e_3}(i, t)}{\sqrt{3}} S_i^{e_3}$$



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Caveat: U(1) is absent

Define spin transport operator:

$$\partial_t S_i^{e_3} + \nabla \cdot \vec{\mathcal{J}} = \mathcal{G}$$

(*For Heisenberg interaction $\mathcal{G} = 0$)

transport operator along \hat{n}_1 , \hat{n}_2 are:

$$\mathcal{J}_1(i) \propto (S_i^y - S_i^z) S_{i+x}^x$$

$$\mathcal{J}_2(i) \propto (S_i^z - S_i^x) S_{i+y}^y$$

Spin Conductivity

For long wave length:

$$\langle \mathcal{J}_n(\mathbf{q}, \omega) \rangle = \underbrace{-\frac{\langle -\mathcal{K}_n^0 \rangle - \Gamma_n(\mathbf{q}, \omega)}{i(\omega + i0^+)} h^{e_3}(\mathbf{q}, \omega)}_{\text{"Drift" } \Sigma} - \underbrace{\frac{\langle -\mathcal{K}_n^1 \rangle - \Lambda_{nm}(\mathbf{q}, \omega)}{i(\omega + i0^+)} iq_m h^{e_3}(\mathbf{q}, \omega)}_{\text{"Diffusive" } \sigma}$$

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$$\Gamma_n = i \int_0^\infty e^{i\omega t} \langle [\mathcal{J}_n(\mathbf{q}, t), \mathcal{G}(-\mathbf{q}, 0)] \rangle$$

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$$\langle \mathcal{K}_n^0 \rangle = -\frac{2}{\sqrt{3}} \langle S_i^n S_{i+n}^n \rangle \sim \langle c_i c_{i+n} \rangle$$

$$\langle \mathcal{K}_n^1 \rangle = 0$$

“Diffusive” Conductivity

- Spin- $\frac{1}{2}$ Heisenberg model

Only diffusive conductivity:

$$\sigma_{nm}(q, \omega) = -\frac{\langle -\mathcal{K} \rangle - \Lambda_{nm}}{i(\omega + i0^+)}$$

In the Drude form

$$\sigma_{nm}(\omega) = D_{\text{diff}}(\omega)\delta(\omega) + \sigma_{nm}^{\text{reg}}(\omega)$$

$$D_{\text{diff}} = 0$$

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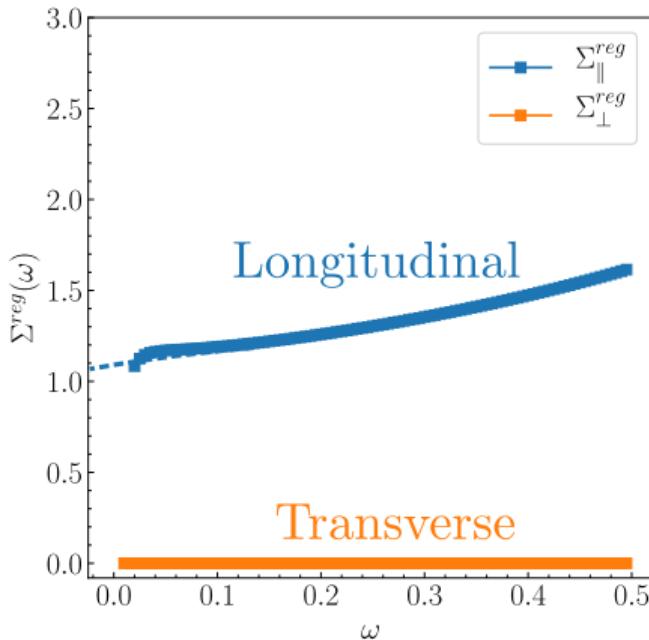
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12-site ED. See also M. Sentef, M. Kollar, and A. P. Kampf, Phys. Rev. B 75, 214403 (2007)

“Diffusive” Conductivity

- Kitaev model

For any (K_x, K_y, K_z) :

$$\langle -\mathcal{K}_n^1 \rangle = 0$$

$$\sigma_{nm}(\mathbf{q}, \omega) = \frac{\Lambda_{nm}}{i(\omega + i0^+)}$$

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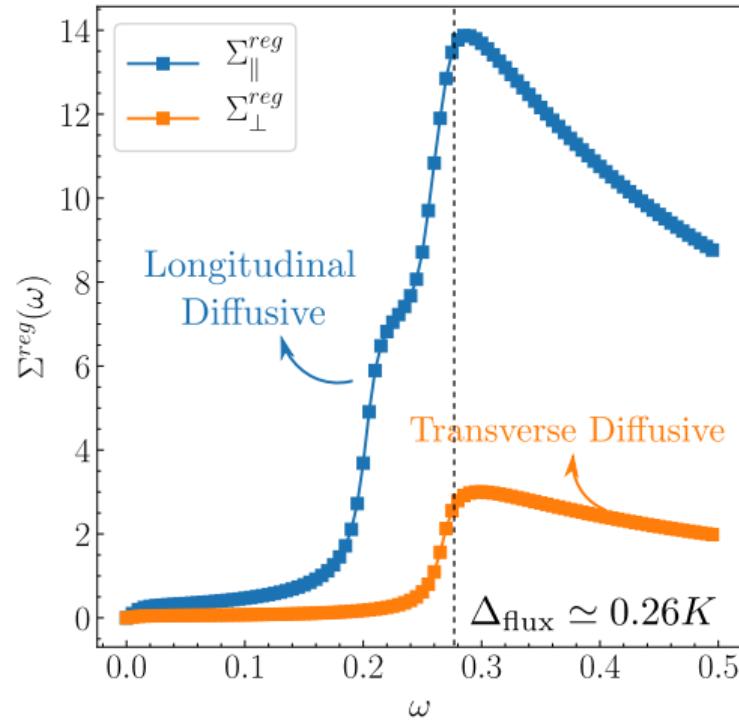
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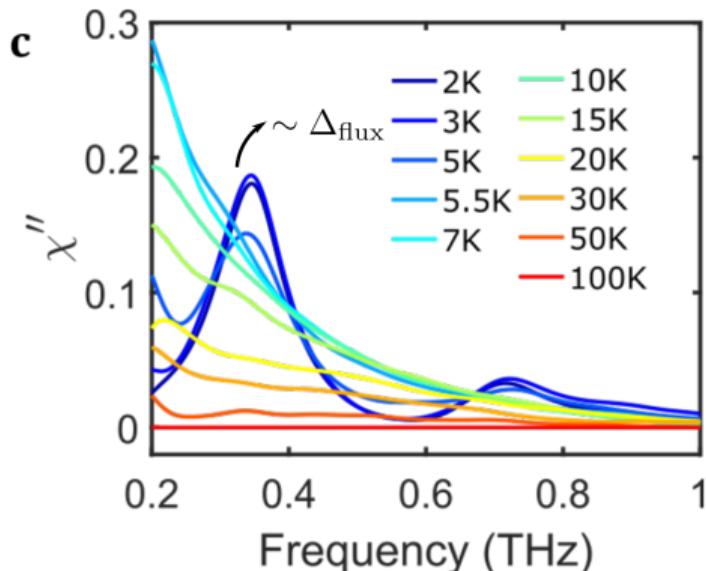
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Magnetic Absorption



Xinshu Zhang, Yuanyuan Xu, Ruidan Zhong, R. J. Cava, N. Drichko, N. P. Armitage,
arXiv:2106.13418 (2021)

“Drift” Conductivity

The "Drift" Conductivity:

$$\Sigma_n(\mathbf{q}, \omega) = -\frac{\langle -\mathcal{K}_n^0 \rangle - \Gamma_n(\mathbf{q}, \omega)}{i(\omega + i0^+)}$$

In the Drude form:

$$\Sigma_n(\omega) = D_{\text{drift}}\delta(\omega) + \Sigma_n^{\text{reg}}(\omega)$$

$$\Sigma_n^{\text{reg}}(\omega) \approx 0$$

with spin Drude weight D :

$$D_{\text{drift}} \approx \langle -\mathcal{K}_n^0 \rangle = \langle c_i c_j \rangle \neq 0$$

unless one of $K_\alpha \rightarrow \infty$

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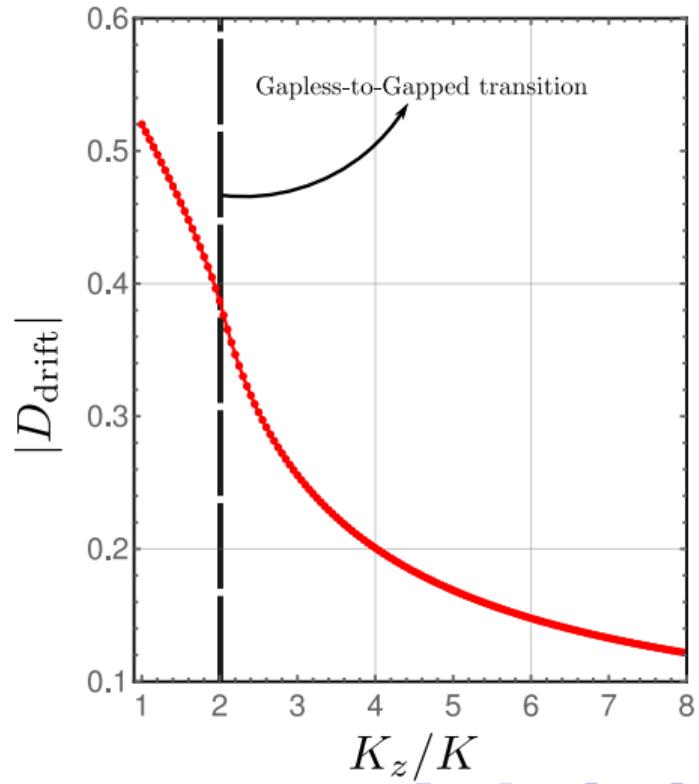
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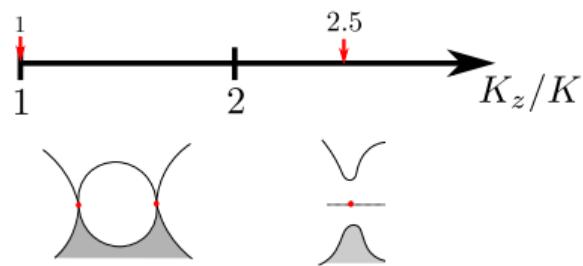
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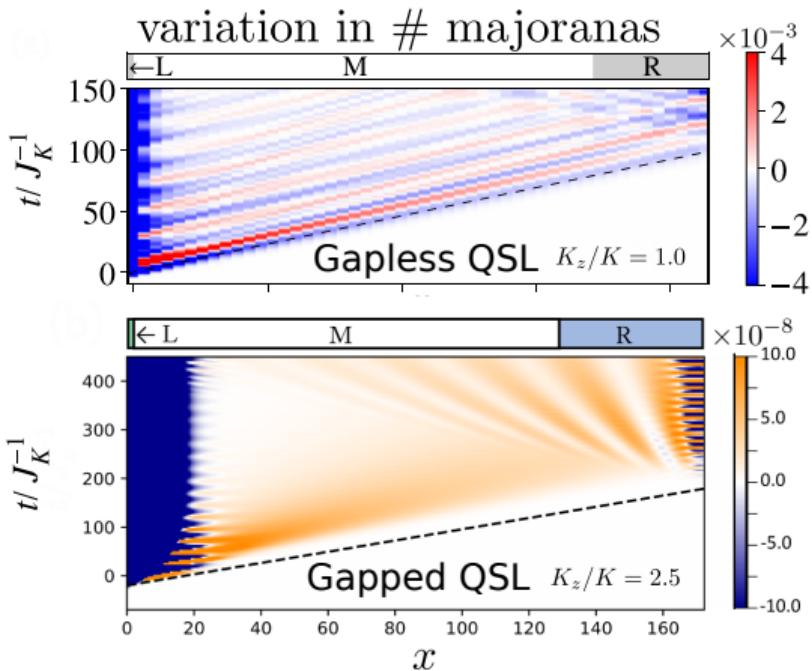


“Drift” Conductivity



Majoranas are able to propagate in both

- ① gapless QSL ($K_z/K = 1.0$)
- ② gapped QSL ($K_z/K = 2.5$)



Minakawa *et al*, PRL 125, 047204 (2020); Taguchi *et al*, PRB 104, 125139 (2021)

Summary

We studied the linear order spin transport of Kitaev honeycomb model in response to a magnetic field in [111] direction, and found that

- ① Effective spin current conductivity can be decomposed into:
 - ① Drift spin conductivity Σ
 - ② Diffusive spin conductivity σ
- ② σ : $\sigma^{\text{reg}} \neq 0$, $D_{\text{diff}} = 0$, the former is responsible for magnetic absorption
- ③ Σ : $\Sigma^{\text{reg}} = 0$, $D_{\text{drift}} \neq 0$, the latter gives ballistic spin transport even in the gapped QSL